

272(2) : Transition to Precessing Planar Ellipse

The beta ellipse is :

$$r = \frac{d}{1 + \epsilon \cos \beta} \quad \text{--- (1)}$$

w/ ϕ $\tan \beta = \frac{L_2}{L} \tan \phi$ --- (2)

Therefore: $\tan \beta = \frac{\sin \beta}{\cos \beta} = \frac{(1 - \cos^2 \beta)^{1/2}}{\cos \beta} = \frac{L_2}{L} \tan \phi$ --- (3)

i.e. $\cos^2 \beta = \frac{1}{1 + \left(\frac{L_2}{L}\right)^2 \tan^2 \phi}$ --- (4)

$$= \frac{\cos^2 \phi}{\cos^2 \phi + \left(\frac{L_2}{L}\right)^2 \sin^2 \phi}$$

It is clear that as :

$$L \rightarrow L_2 \quad \text{--- (5)}$$

$$\beta \rightarrow \phi \quad \text{--- (6)}$$

Now define :

$$\cos(x\phi) = \frac{\cos \phi}{\left(\cos^2 \phi + \left(\frac{L_2}{L}\right)^2 \sin^2 \phi\right)^{1/2}} \quad \text{--- (7)}$$

so $x = \frac{1}{\phi} \cos^{-1} \left(\frac{\cos \phi}{\left(\cos^2 \phi + \left(\frac{L_2}{L}\right)^2 \sin^2 \phi\right)^{1/2}} \right)$ --- (8)

2) In general therefore the 3-D beam gives the processing ellipse:

$$r = \frac{d}{1 + \epsilon \cos(x\phi)} \quad - (9)$$

with x defined by Eq. (8).

It is asserted experimentally that:

$$x = 1 + \frac{3MG}{dc^2} \quad - (10)$$

so eqs. (8) and (10) determine an angle ϕ_0 :

$$\frac{1}{\phi_0} \cos^{-1} \left(\frac{\cos \phi_0}{\left(\cos^2 \phi_0 + \left(\frac{Lz}{L} \right)^2 \sin^2 \phi_0 \right)^{1/2}} \right) = 1 + \frac{3MG}{dc^2}$$

hence three dimensional orbits have evolved to processing two dimensional orbits, which process.

In general however:

$$r = \frac{d}{1 + \epsilon \cos \beta} \quad - (12)$$

where

$$\cos \beta = \frac{\cos \phi}{\left(\cos^2 \phi + \left(\frac{Lz}{L} \right)^2 \sin^2 \phi \right)^{1/2}} \quad - (13)$$

It is clear that eq. (12) is no longer an

ellipse. Therefore primordial three dimensional orbits evolve into the desired two dimensional precessing orbits.
 The time dependence of this process is given by:

$$\frac{d\beta}{dt} = \frac{L}{mr^2} \quad - (14)$$

So
$$dt = \frac{mr^2}{L} d\beta \quad - (15)$$

and
$$t = \frac{md^2}{L} \int \frac{d\beta}{(1 + \epsilon \cos \beta)^2} \quad - (16)$$

So t can be evaluated in terms of β and ϕ .
 To evaluate eq. (16) we use Kepler's equation

$$\frac{2\pi t}{\tau} = \phi - \epsilon \sin \phi \quad - (17)$$

for the near anomaly, where:

$$\tan \frac{\beta}{2} = \left(\frac{1+\epsilon}{1-\epsilon} \right)^{1/2} \tan \frac{\phi}{2} \quad - (18)$$

Here τ is the time taken to sweep out an entire orbit.

Now use the half angle formula:

$$\tan^2 \frac{A}{2} = \frac{1 - \cos A}{1 + \cos A} \quad - (19)$$

So:
$$\frac{1 - \cos \beta}{1 + \cos \beta} = \left(\frac{1 + \epsilon}{1 - \epsilon} \right) \left(\frac{1 - \cos \phi}{1 + \cos \phi} \right) \quad - (20)$$

where:
$$\cos \beta = \frac{\cos \phi}{\left(\cos^2 \phi + \left(\frac{L_2}{L} \right)^2 \sin^2 \phi \right)^{1/2}} \quad - (21)$$

so ϕ and $\sin \phi$ can be found in terms of $\cos \phi$, and t evaluated.

The quantity $\cos \beta$ can also be expressed in terms of θ :

$$\cos^2 \beta = 1 - \frac{L^2}{L^2 - L_2^2} \cos^2 \theta \quad - (22)$$

i.e.
$$(L^2 - L_2^2) \cos^2 \beta = L^2 (1 - \cos^2 \theta) - L_2^2 \quad - (23)$$

$$= (L^2 - L_2^2) \frac{\cos^2 \phi}{\cos^2 \phi + \left(\frac{L_2}{L} \right)^2 \sin^2 \phi}$$

It is clear that θ is a function of ϕ and can be expressed in terms of ϕ and is due to the use of a true function. Eq. (23) removes any restriction.

5) Suggested Graphics

1) Plot the function:

$$r = \frac{d}{1 + \epsilon \cos \phi} \quad - (24)$$

and compare it directly with:

$$r = \frac{d}{1 + \epsilon \cos \phi} \quad - (25)$$

2) Plot the function:

$$x = \frac{1}{\phi} \cos^{-1} \left(\frac{\cos \phi}{\left(\cos^2 \phi + \left(\frac{L_2}{L} \right)^2 \sin^2 \phi \right)^{1/2}} \right) \quad - (26)$$

3) Plot and animate t versus ϕ for eq. (16).
