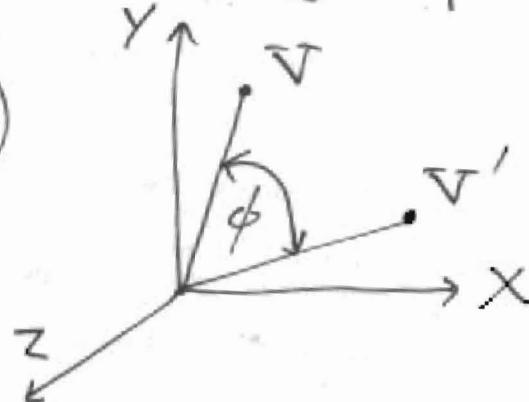


1) 153(12): Conversion Matrix and Rotation Generator.

It was shown in note 153(11) that the tetrad matrix is the transformation matrix. Consider the rotation of a vector in 3-D in the xz plane about the z axis (Ryder p. 30):

$$\begin{bmatrix} v_x' \\ v_y' \\ v_z' \end{bmatrix} = \begin{bmatrix} \cos \phi & \sin \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} v_x \\ v_y \\ v_z \end{bmatrix} \quad (1)$$



The passive rotation is defined as the rotation of a vector about the z axis by rotating the axes anticlockwise (see Ryder p. 30). The rotation of the axes introduce a convention. The rotation generator is defined by:

$$J_z = \frac{1}{i} \left. \frac{dR_z(\phi)}{d\phi} \right|_{\phi=0} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad (2)$$

where

$$R_z(\phi) = \begin{bmatrix} \cos \phi & \sin \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (3)$$

so, s in Ryder:

$$[J_x, J_y] = i J_z \quad (4)$$

etc cyclicum

The regular momentum operators of quantum mechanics are

$$[J_x, J_y] = i \hbar J_z \quad (5)$$

etc cyclicum.

These operators originate in the tetrad:

2)

$$\underline{v}^a = \begin{bmatrix} \cos\phi & \sin\phi & 0 \\ -\sin\phi & \cos\phi & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} v_1^{(1)} & v_2^{(1)} & 0 \\ v_1^{(2)} & v_2^{(2)} & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad -(6)$$

where, is the cylindrical polar system:

$$\underline{e}_r = \underline{e}_1 = \underline{e}^{(1)} = \underline{v}^{(1)} = v_1^{(1)} \underline{i} + v_2^{(1)} \underline{j} \quad -(7)$$

$$\underline{e}_\phi = \underline{e}_2 = \underline{e}^{(2)} = \underline{v}^{(2)} = v_1^{(2)} \underline{i} + v_2^{(2)} \underline{j} \quad -(8)$$

$$\underline{e}_z = \underline{e}_3 = \underline{e}^{(3)} = \underline{v}^{(3)} = \underline{k} \quad -(9)$$

In Cartan geometry this process is governed by the tetrad postulate, because rotation conserves the vector field, so:

$$\begin{aligned} D_\mu v^\alpha &= \partial_\mu v^\alpha + \omega_{\mu b}^\alpha v^b - \Gamma_{\mu\nu}^\lambda v^\lambda = 0 \quad -(10) \\ &= \partial_\mu v^\alpha + \omega_{\mu\nu}^\alpha - \Gamma_{\mu\nu}^\alpha \end{aligned}$$

Therefore

$$\boxed{\partial_\mu v^\alpha = \underline{d}_{\mu\nu}^\alpha} \quad -(11)$$

where:

$$\underline{d}_{\mu\nu}^\alpha = \underline{\Gamma}_{\mu\nu}^\alpha - \underline{\omega}_{\mu\nu}^\alpha \quad -(12)$$

So the conservation matrix $\underline{d}_{\mu\nu}^\alpha$ is a rotation generator matrix and an angular momentum operator.

In the notation of eq. (6) the lower indices without brackets are indices of the Cartesian system and the upper indices with brackets are indices of the cylindrical polar system.

3) This system is defined by:

$$\cos \phi = \frac{x}{r}, \quad \sin \phi = \frac{y}{r}, \quad z = z \quad -(13)$$

where

$$r = (x^2 + y^2)^{1/2} \quad -(14)$$

= constant

Therefore: $\frac{\partial \sqrt{r}}{\partial y} = \frac{\partial \sqrt{r}}{\partial y} = \frac{1}{r} \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad -(15)$

$$\frac{\partial \sqrt{r}}{\partial x} = \frac{\partial \sqrt{r}}{\partial x} = \frac{1}{r} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad -(16)$$

So:

$$\frac{1}{r} \begin{bmatrix} 0 & 1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & d_{22}^{(1)} & 0 \\ d_{21}^{(2)} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} \quad -(17)$$

i.e.

$d_{22}^{(1)} = \frac{1}{r}, \quad d_{21}^{(2)} = -\frac{1}{r}$

$$-(18)$$

As in paper 63, the conversion are inversely proportional
to the radial component r .

$$\begin{bmatrix} 0 & d_{22}^{(1)} & 0 \\ d_{21}^{(2)} & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix} = \frac{i}{r} J_z \quad -(19)$$

Q.E.D.

The conversion matrix is the rotation generator
matrix with i/r , and the angular momentum operator

4) matrix with rk : / (l_r)

These calculations have been built up from consideration only of the cylindrical polar coordinate system and how it is related to the Cartesian coordinate system. The whole of quantum mechanics in molecular spectroscopy (see Atkins) depends on the angular momentum operators.

The metric is defined by :

$$g_{\mu\nu} = g^a_\mu g^b_\nu \eta - (20)$$

so $g_{11} = g^{(1)}_1 g^{(1)}_1 \eta_{(1)(1)} + g^{(2)}_1 g^{(2)}_1 \eta_{(2)(2)} - (21)$
 $= \cos^2 \phi + \sin^2 \phi = 1$

if $\eta_{(1)(1)} = \eta_{(2)(2)} = 1 - (22)$

so:
$$\boxed{g_{\mu\nu} = \eta_{ab} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}} - (23)$$

Cartesian
Torsion is defined by :

$$\begin{aligned} T^a_{\mu\nu} &= \partial_\mu g^a_\nu - \partial_\nu g^a_\mu + \omega^a_{\mu\nu} - \omega^a_{\nu\mu} \\ &= \Gamma^a_{\mu\nu} - \Gamma^a_{\nu\mu} \end{aligned} - (24)$$

and there are non-zero elements such as :

$$\begin{aligned} T^{(2)}_{12} &= \partial_1 g^{(2)}_2 - \partial_2 g^{(2)}_1 + \omega^{(2)}_{12} - \omega^{(2)}_{21} \\ &= 2/r + \omega^{(2)}_{12} - \omega^{(2)}_{21} \end{aligned} - (25)$$

5) The Riemann torsion is defined by:

$$T_{\mu\nu}^{\lambda} = \nabla_{\mu}^{\lambda} T_{\nu\nu} - (26)$$

so elements of the Riemann torsion also exist. The Riemann torsion is:

$$T_{\mu\nu}^{\lambda} = \Gamma_{\mu\nu}^{\lambda} - \Gamma_{\nu\mu}^{\lambda} - (27)$$

so the $\Gamma_{\mu\nu}^{\lambda}$ covariant also exists. The latter is defined by:

$$[D_{\mu}, D_{\nu}] \nabla^{\sigma} = R^{\sigma}_{\lambda\mu\nu} \nabla^{\lambda} - T^{\rho}_{\mu\nu} D_{\rho} \nabla^{\sigma} - (28)$$

$$= - \Gamma_{\mu\nu}^{\rho} D_{\rho} \nabla^{\sigma} +$$

$$\boxed{\Gamma_{\mu\nu}^{\lambda} = - \Gamma_{\nu\mu}^{\lambda}} - (29)$$

so:

$$[D_{\mu}, D_{\nu}] = - [D_{\nu}, D_{\mu}] - (30)$$

The covariant derivative is:

$$D_{\mu} X^{\nu} = \partial_{\mu} X^{\nu} + \tilde{\Gamma}_{\mu\lambda}^{\nu} X^{\lambda} - (31)$$

so is different from $\partial_{\mu} X^{\nu}$ in general. Thus difference is due to the passive rotation of eq.(1).

6) Finally, eq. (11) is the mathematical expression of the fact that a passive rotation is equivalent to the usual way in which a rotation is thought of - as a passive rotation of a vector keeping coords fixed. The active rotation is d_{μ}^{ν} and the passive rotation is $d_{\mu\nu}^{\nu}$.

Obviously, in the usual way of dealing with rotations, the idea of conservation is never considered in what is referred to as "three dimensional space", but the conservation is inherent in the analysis. The very definition of the cylindrical polar coords is enough to generate a conservation.
