

260(7) : Development of the Subsidiary Condition

As shown in note 260(6) the subsidiary condition is

$$\underline{p} \cdot \underline{\nabla} (\underline{\nabla} \cdot \underline{p}) = 0 \quad - (1)$$

Under this condition the Beltrami momentum equation

$$\underline{\nabla} \times \underline{p} = \kappa \underline{p} \quad - (2)$$

gives the Schrödinger equation:

$$(\nabla^2 + \kappa^2) \psi = 0 \quad - (3)$$

where

$$\kappa^2 = \frac{2m}{\hbar^2} (V - E) \quad - (4)$$

Two possible solutions of eq. (1) are:

$$\underline{\nabla} \cdot \underline{p} = 0 \quad - (5)$$

and

$$\underline{\nabla} (\underline{\nabla} \cdot \underline{p}) = 0 \quad - (6)$$

with

$$\underline{p} = -i\hbar \underline{\nabla} \quad - (7)$$

Eq. (5) gives:

$$\nabla^2 \psi = 0 \quad - (8)$$

which is consistent with eq. (3) only if

$$\kappa = 0 \quad - (8)$$

2) Eq. (6) give:

$$\underline{\nabla} (\nabla^2 \phi) = 0 \quad - (9)$$

where $\nabla^2 \phi = -\kappa^2 \phi \quad - (10)$

so $\underline{\nabla} (\kappa^2 \phi) = 0 \quad - (11)$
 $= (\underline{\nabla} \kappa^2) \phi + \kappa^2 \underline{\nabla} \phi$

Therefore: $\boxed{\underline{\nabla} \phi = - \left(\frac{\underline{\nabla} \kappa^2}{\kappa^2} \right) \phi} \quad - (12)$

and $\underline{\nabla} \cdot \underline{\nabla} \phi = \nabla^2 \phi = -\underline{\nabla} \cdot \left(\frac{\underline{\nabla} \kappa^2}{\kappa^2} \phi \right)$
 $= - \left(\left(\underline{\nabla} \cdot \left(\frac{\underline{\nabla} \kappa^2}{\kappa^2} \right) \right) \phi + \left(\frac{\underline{\nabla} \kappa^2}{\kappa^2} \right) \cdot \underline{\nabla} \phi \right) \quad - (13)$
 $= - \left(\underline{\nabla} \cdot \frac{\underline{\nabla} \kappa^2}{\kappa^2} + \frac{\underline{\nabla} \kappa^2}{\kappa^2} \cdot \frac{\underline{\nabla} \kappa^2}{\kappa^2} \right) \phi$

From a comparison of eqns. (3) and (13)
we obtain the subsidiary condition:

$$\kappa^2 = \frac{\nabla \cdot \nabla \kappa^2}{\kappa^2} = \frac{\nabla \kappa^2 \cdot \nabla \kappa^2}{\kappa^4} \quad - (14)$$

$$= \frac{\nabla^2 \kappa^2}{\kappa^2} - \frac{\nabla \kappa^2 \cdot \nabla \kappa^2}{\kappa^4}$$

i.e. $\boxed{\kappa^2 \nabla^2 \kappa^2 = \nabla \kappa^2 \cdot \nabla \kappa^2 + \kappa^6}$ - (15)

where

$$\kappa^2 = \frac{2m}{\hbar^2} (V - E) \quad - (16)$$

Eqs. (15) and (16) define the potential V needed for the momentum Beltrami equation to become a Schrodinger equation.

Finally, this potential is used to calculate the scattering amplitude in the Born approximation:

$$f(\hat{\underline{\kappa}}') = -\frac{m}{2\pi\hbar^2} \int e^{-i\underline{\kappa} \cdot \underline{r}'} V(\underline{r}') e^{i\underline{\kappa} \cdot \underline{r}'} d\tau' \quad - (17)$$